

Liouville $\rightarrow$ BBGKY

Born

Bohrlubov

Green

Kirkwood

von

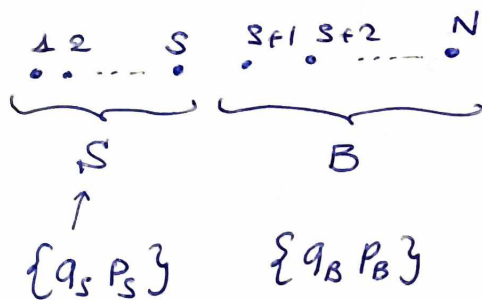
• Liouville's theorem

$$\frac{\partial \rho}{\partial t} = -\{\rho, H\} = \frac{\partial H}{\partial q} \frac{\partial \rho}{\partial p} - \frac{\partial H}{\partial p} \frac{\partial \rho}{\partial q}$$

Vs Quantum

$$\frac{\partial \hat{\rho}}{\partial t} = \frac{1}{i\hbar} [\hat{H}, \hat{\rho}]$$

• Partition



$$H = H_S(q_S, p_S) + H_B(q_B, p_B) + H_{SB}(q_S, q_B, p_S, p_B)$$

• Goal

from $\rho^{\text{TOT}}(q, p, t)$ determine $\rho^{(S)}(q, \dots, q_S, p, \dots, p_S, t)$

classical system

$$\rho^{(s)} = \text{tr}_B \rho^{\text{TOT}}$$

$$\rho^{(s)} = \int \prod_{k=1}^N d^3 p_k d^3 q_k \rho^{\text{TOT}}$$

$$\text{tr}_B \Leftrightarrow \int \prod_{k=1}^N d^3 p_k d^3 q_k$$

classical system

$$\frac{\partial \rho^{(s)}}{\partial t} = \text{tr}_B \{H_S, \rho\} +$$

~~$$\text{tr}_B \{H_B, \rho\} +$$~~

$$\text{tr}_B \{H_{SB}, \rho\}$$

→ generalised Gauss theorem

$$\int d^d x \nabla \cdot \vec{a}(x) = - \int d^d x \rho(x) \nabla \cdot \vec{a}(x)$$

$$\int d^d p_B d^d q_B \left[\underbrace{\frac{\partial H_B}{\partial q_B} \frac{\partial \rho}{\partial p_B}}_{(a)} - \underbrace{\frac{\partial H}{\partial p_B} \frac{\partial \rho}{\partial q_B}}_{(b)} \right]$$

a) integrating in $d p_B$

$$- \int d^d q_B d^d p_B \frac{\partial^2 H}{\partial q_B \partial p_B} \cdot \rho$$

b) integrating in $d q_B$

$$+ \int d^d p_B d^d q_B \frac{\partial^2 H}{\partial p_B \partial q_B} \rho$$

the difference vanishes

Now consider the third term

$$\int dq_B dp_B \left\{ \frac{\partial H_{SB}}{\partial q_B} \frac{\partial p}{\partial p_B} - \frac{\partial H_{SB}}{\partial p_B} \frac{\partial p}{\partial q_B} + \frac{\partial H_{SB}}{\partial q_S} \frac{\partial p}{\partial q_S} - \frac{\partial H_{SB}}{\partial p_S} \frac{\partial p}{\partial q_S} \right\}$$

↑
depends only on coordinates

• Pairwise interaction

$$H = \sum_i \frac{p_i^2}{2m} + \frac{1}{2} \sum_{i \neq j} V(|\vec{q}_i - \vec{q}_j|) + \sum_i \underbrace{\phi(q_i)}_{\text{ext potential}}$$

therefore

$$H_{SB} = \sum_i^S \sum_{j=S+1}^N V(|q_i^S - q_j^B|) \quad (\text{no factor } \frac{1}{2} !!)$$

the first term reads

$$\int \prod_{k=S+1}^N \frac{d^3 p_k}{\pi_k} d^3 q_k \sum_i^S \sum_{j=S+1}^N \frac{\partial V}{\partial q_j} \frac{\partial p}{\partial p_j} + \frac{\partial V}{\partial q_i} \frac{\partial p}{\partial p_i}$$

$$= \int \prod_{k=S+1}^N \frac{d^3 p_k}{\pi_k} d^3 q_k \sum_i^S \sum_{j=S+1}^N \hat{\Theta}_{ij} p^{\text{TOT}}$$

$$\Theta_{ij} = \frac{\partial V}{\partial q_j} \frac{\partial}{\partial p_j} + \frac{\partial V}{\partial q_i} \frac{\partial}{\partial p_i}$$

integrated variables on identical particles

$$= \sum_i^S (N-S) \int \prod_{k=S+1}^N \frac{d^3 p_k}{\pi_k} d^3 q_k \Theta_{i, S+1} p^{\text{TOT}}$$

$\underbrace{d^3 p_{S+1} d^3 q_{S+1} \prod_{k=S+2}^N \frac{d^3 p_k}{\pi_k} d^3 q_k}_{\text{integrated}}$

$$\text{tr}_B \{H_{SB}, \rho\} = (N-s) \sum_i^s \int d^3 q_{s+1} d^3 p_{s+1} \hat{\Theta}_{s+1} \rho^{(s+1)}$$

$$\int d^3 p_{s+1} d^3 q_{s+1} \left[\underbrace{\frac{\partial V_{s+1}}{\partial q_\alpha}}_{(a)} \frac{\partial \rho^{(s+1)}}{\partial p_\alpha} + \underbrace{\frac{\partial V_{s+1}}{\partial q_{s+1}} \frac{\partial \rho^{(s+1)}}{\partial p_{s+1}}}_{(b)} \right]$$

(a) (b)
Generalised Gauss

b) integrate in p_{s+1}

$$- \int d^3 q_{s+1} d^3 p_{s+1} \frac{\partial^2 V_{s+1}}{\partial p_{s+1} \partial q_{s+1}} \rho^{(s+1)}$$

Indep. on p

$$\text{tr}_B \{H_{SB}, \rho\} = \int d^3 p_{s+1} d^3 q_{s+1} (N-s) \sum_i^s \left(\frac{\partial V_{s+1}}{\partial q_\alpha} \right) \frac{\partial \rho^{(s+1)}}{\partial p_\alpha}$$

$-\vec{F}_{s+1}$

Finally

$$\frac{\partial \rho^{(s)}}{\partial t} - \{H_s, \rho^{(s)}\} = -(N-s) \sum_i^s \int d^3 p_{s+1} d^3 q_{s+1} \vec{F}_{s+1} \cdot \frac{\partial \rho^{(s+1)}}{\partial \vec{p}_\alpha}$$

Examples

$S=1$ cluster of only 1 particle

$$H_S = \frac{\vec{p}^2}{2m} + \underbrace{\phi(\vec{q})}_{\text{external potential}} = H^{(1)} \quad S_S = \rho^{(1)}(\vec{p}, \vec{q}, t)$$

$$\{H_S, S_S\} = \frac{\partial H}{\partial \vec{q}} \cdot \frac{\partial \rho^{(1)}}{\partial \vec{p}} - \frac{\partial H}{\partial \vec{p}} \cdot \frac{\partial \rho^{(1)}}{\partial \vec{q}} = \frac{\partial \phi}{\partial \vec{q}} \cdot \frac{\partial \rho^{(1)}}{\partial \vec{p}} - \frac{\vec{p}}{m} \cdot \frac{\partial \rho^{(1)}}{\partial \vec{q}}$$

$$\frac{\partial \rho^{(1)}}{\partial t} - \{H_S, \rho^{(1)}\} = -(N-1) \int d\vec{p}_2 d\vec{q}_2 F(\vec{q} - \vec{q}_2) \cdot \frac{\partial \rho^{(2)}(\vec{q}, \vec{q}_2, \vec{p}, \vec{p}_2, t)}{\partial \vec{p}}$$

$$\left\{ \frac{\partial}{\partial t} + \frac{\vec{p}}{m} \cdot \frac{\partial}{\partial \vec{q}} - \frac{\partial \phi(\vec{q})}{\partial \vec{q}} \cdot \frac{\partial}{\partial \vec{p}} \right\} \rho^{(1)}(\vec{q}, \vec{p}, t) = -(N-1) \int d\vec{p}_2 \int d\vec{q}_2 F(\vec{q} - \vec{q}_2) \frac{\partial \rho^{(2)}}{\partial \vec{p}}$$

$S=2$

$$H_S = H^{(2)} = \frac{\vec{p}_1^2}{2m} + \frac{\vec{p}_2^2}{2m} + \phi(\vec{q}_1) + \phi(\vec{q}_2) + V(\vec{q}_1 - \vec{q}_2)$$

$$\left\{ \frac{\partial}{\partial t} + \sum_{i=1}^2 \frac{\vec{p}_i}{m} \cdot \frac{\partial}{\partial \vec{q}_i} - \frac{\partial \phi(\vec{q}_1)}{\partial \vec{q}_1} \cdot \frac{\partial}{\partial \vec{p}_1} - \Theta_{12} \right\} \rho^{(2)}(1, 2, t) =$$

$$= + (N-2) \int d\vec{p}_3 d\vec{q}_3 \{ \hat{\Theta}_{13} + \hat{\Theta}_{23} \} \rho^{(3)}(1, 2, 3, t)$$

$$\Theta_{ij} = \frac{\partial V}{\partial \vec{q}_j} \cdot \frac{\partial}{\partial \vec{p}_i} + \frac{\partial V}{\partial \vec{q}_i} \cdot \frac{\partial}{\partial \vec{p}_j}$$