

Ginsburg Criterion for the validity of the mean-field theory

①

Mean field theory neglects correlations that is is equivalent to consider an effective Hamiltonian in which the quadratic interaction is linearised

Another way of introducing MFT in Ising is to consider each spin as

$$\sigma_i = \langle \sigma_i \rangle + \delta \sigma_i \quad (\text{formal } \sigma_i = \pm 1)$$

$$H[\sigma] = -\frac{J}{2} \sum_{\langle ij \rangle} \sigma_i \sigma_j - h \sum_i \sigma_i$$

$$H[\sigma] = H[\langle \sigma \rangle] + \underbrace{\frac{\partial H}{\partial \sigma}}_{\text{functional derivative}} \Big|_{\langle \sigma \rangle} \delta \sigma + \mathcal{O}(\delta \sigma^2)$$

$$H[\sigma] \approx H[\langle \sigma \rangle] - J \sum_i \delta \sigma_i \sum_j^{(i)} \langle \sigma_j \rangle - h \sum_i \delta \sigma_i$$

the replacing again $\delta \sigma_i = \sigma_i - \langle \sigma_i \rangle$

one gets with $m_i = \langle \sigma_i \rangle$

$$H \approx -\frac{J}{2} \sum_{\langle ij \rangle} m_i m_j - h \sum_i m_i$$

subtract

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$$-J \sum_i (\sigma_i - m_i) \sum_j^{(i)} m_j - h \sum_i (\sigma_i - m_i)$$

$$= \frac{J}{2} \sum_{\langle ij \rangle} m_i m_j - \sum_i \sigma_i (\sum_j^{(i)} m_j + h) = H_{\text{eff}}$$

and

$$\frac{1}{3} \ln \mathcal{Z}(m) = \text{tr}_{\{\sigma\}} e^{\beta H_{\text{eff}}}$$

with

$$\mathcal{Z}(m) = \frac{J}{2} \sum_{\langle ij \rangle} m_i m_j - \frac{1}{\beta} \sum_i \ln 2 \cosh \beta [J \sum_j^{(i)} m_j + h]$$

Landau potential

$$(\langle \sigma \rangle) \sigma + \Delta \sigma \frac{\partial \mathcal{L}}{\partial H} + (\langle \sigma \rangle) H = H$$

$$\sigma = \langle \sigma \rangle + \Delta \sigma$$

$$H(\sigma) = -\frac{J}{2} \sum_{\langle ij \rangle} \sigma_i \sigma_j - h \sum_i \sigma_i$$

therefore MFT hold when the correlation function (3) really factorises

$$C(\vec{R}) = \beta^{-1} \chi(\vec{R}) \quad \text{from Linear Response}$$

$$\left. \begin{aligned} \frac{C(\vec{R})}{m^2} &\ll 1 \\ \frac{\langle \sigma_{i+R} \sigma_i \rangle}{\langle \sigma_{i+R} \rangle \langle \sigma_i \rangle} &\ll 1 \end{aligned} \right\} \begin{array}{l} \text{Ginzburg} \\ \text{Criterion} \\ (T > T_c) \end{array}$$

evaluate $C(\vec{R})$ as $\beta^{-1} \chi(\vec{R})$

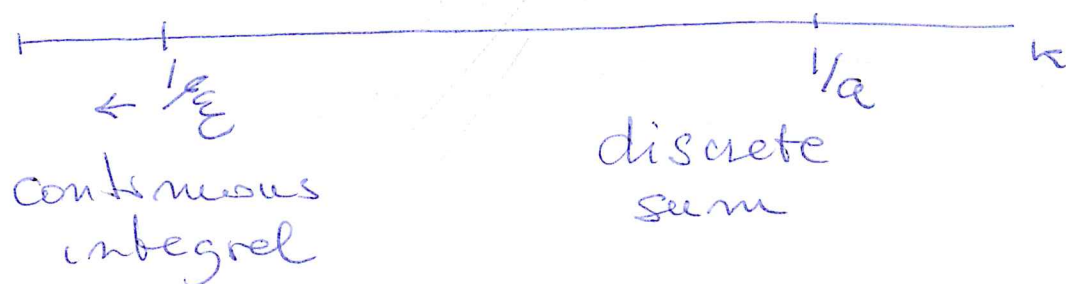
and

$$\chi(\vec{R}) = \frac{1}{N} \sum_{\vec{k}} e^{i\vec{k} \cdot \vec{R}} \chi(\vec{k})$$

$$= \underbrace{\frac{1}{N} \sum_{\vec{k}} e^{i\vec{k} \cdot \vec{R}} \chi(\vec{k})}_{\substack{\text{discrete} \\ \text{regular} \\ \wedge \text{ some infrared} \\ \text{cutoff}}} + \underbrace{\frac{1}{N} \sum_{\vec{k}} e^{i\vec{k} \cdot \vec{R}} \chi(\vec{k})}_{\substack{\text{continuous} \\ \text{divergent at } T_c \\ \downarrow \\ \chi^d}}$$

what is Λ

(4)



assume that $\Lambda \sim 1/\epsilon_c$ such as

$$\chi(\vec{R}) = \chi^d + \frac{1}{N} \sum_{\vec{k} \neq 0} e^{i\vec{k} \cdot \vec{R}} \chi(\vec{k})$$

in the J_{xy} model at MF level

$$\chi(\vec{k}) = \beta \frac{1 - m^2}{1 - (1 - m^2) \beta \mathcal{F}(\vec{k})}$$

with $\mathcal{F}(\vec{k}) = z \mathcal{F} \sum_{\alpha}^d \cos(k_{\alpha} a)$

as $\vec{k} \neq 0$ and $T \rightarrow T_c^-$ $m \rightarrow 0$ and

$$\chi(\vec{k} \neq 0) \rightarrow \beta_c \frac{1}{1 - \beta_c \mathcal{F} z \sum_{\alpha} \cos(k_{\alpha} a)}$$

since $\beta_c \mathcal{F} z = 1$ but $\cos(k_{\alpha} a) < 1$ as $k \neq 0$

the expression do not diverge!

$$\chi^d = \frac{L^d}{(2\pi)^d} \frac{1}{N} \int_0^{\Lambda} d^d k \left(e^{i\vec{k} \cdot \vec{R}} \right) \frac{\chi_0}{1 + (\epsilon_c k)^2} = \frac{a^d}{(2\pi)^d} \Omega_d \int_0^{\Lambda} dk \frac{k^{d-1}}{1 + (\epsilon_c k)^2}$$

approximate
 $k \approx 0$

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$$\chi_{d\text{d}} \sim \frac{d \chi_0}{(2\pi)^d} \Omega_d \int_0^{1/\xi} dk \frac{k^{d-1}}{1 + (\xi^2 k^2)}$$

$$\frac{1}{\xi^d} \int_0^1 dx \frac{x^{d-1}}{1+x^2}$$

finite for any $d > 0$

express all quantities
as a function of ξ

$$\chi_0 \sim \xi^2 \quad \text{so that}$$

$$\chi_{d\text{d}} \sim \frac{1}{\xi^{d-2}} \quad \text{as } \xi \rightarrow \infty \quad \text{diverges when } d < 2$$

$$\begin{aligned} \xi &\sim (T - T_c)^{-\nu} & \nu = 1/2 \\ \chi_0 &\sim (T - T_c)^{-\gamma} & \gamma = 1 \\ m &\sim (T_c - T)^\beta & \beta = 1/2 \end{aligned} \quad \text{MFT exponents}$$

Now compare

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$$\frac{C(\vec{R})}{m^2} = \frac{\chi(\vec{R})}{\beta m^2} \approx \frac{\chi^d(R)}{\beta m^2}$$

Since $m^2 \sim \frac{1}{\xi_a^2}$ we have

$$\left. \frac{C(R)}{m^2} \right\} \approx \frac{\xi_a^2}{\xi_a^{d-2}} = \xi_a^{4-d} \quad \underline{\text{diverges as } d < 4}$$

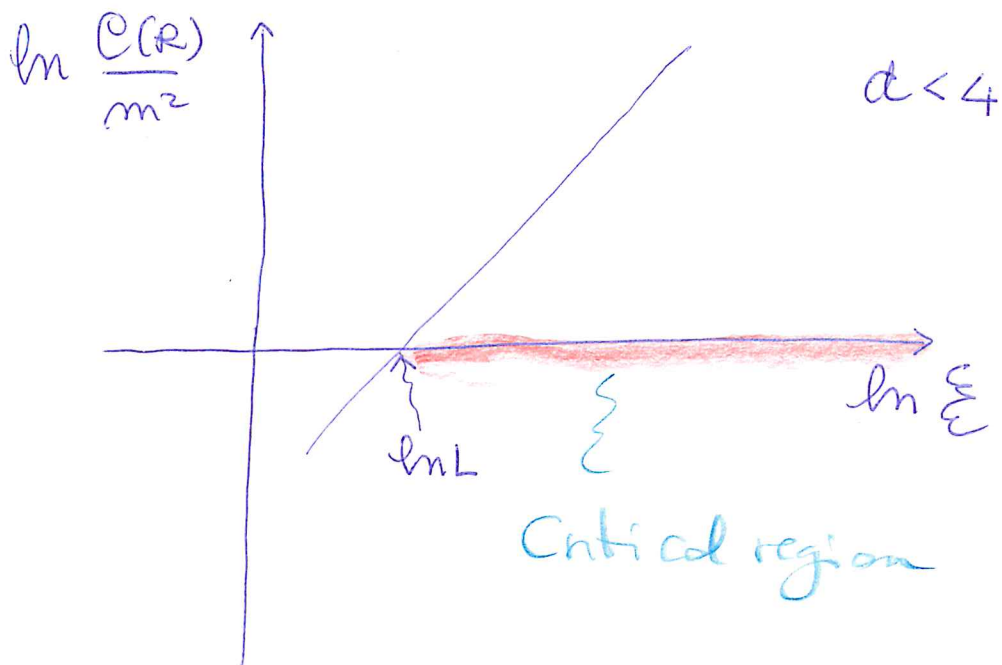
long distances

for dimensional reason

dimensional

$$\left(\frac{C(R)}{m^2} \right) = \left(\frac{\xi_a}{L} \right)^{4-d}$$

where L is a non-universal length which depends on the system's params



$$\ln \frac{C(R)}{m^2} = (4-d) \ln \xi - (4-d) \ln L$$

as $\xi > L \rightarrow$ critical region where
critical exponents
deviates from that
of mean field